

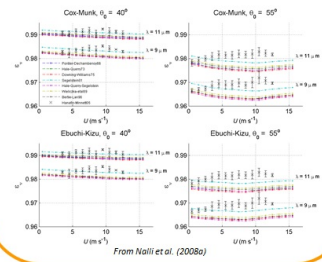
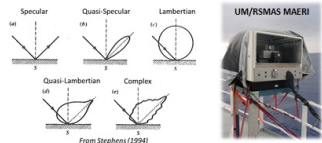
Infrared Sea Surface Effective-Emissivity (IRSSE) Model Upgrade Plans for CRTM

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Background

- For satellite IR remote sensing applications, the **surface emissivity/reflectance** spectrum must be specified with a high degree of absolute accuracy
 - ~0.5% uncertainty results ~0.3-0.4 K systematic error in LWIR window channels
- IR sea-surface emissivity models have gained widespread acceptance (e.g., Masuda et al. 1988; Watts et al. 1996; Wu and Smith 1997), but only after they were validated
 - Masuda's model was published in 1988, but no one used it because it was never validated against observations
 - Marine Atmospheric Emitted Radiance Interferometer (MAERI) (Smith et al. 1996; Minnett et al. 2001) led to acceptance and application of emissivity models
 - These models calculated emissivity as the ensemble-mean of 1 - ρ of surface wave facets
 - The models were improved, but residual systematic discrepancies (0.1-0.4 K) remained at higher wind speeds and view angles 240° (Nalli et al. 2001, 2006; Hanafin and Minnett 2005) due to incorrect specification of reflected atmospheric radiance



JCSDA, STAR and JPSS Support for IRSSE Model Development

- JCSDA and STAR supported in-house FY05-06 research to find a workable solution for application to the CRTM
- This research culminated in the CRTM IRSSE model (Nalli et al. 2018a,b; van Delst et al. 2009)
- Notably, the IRSSE model uses the **effective emissivity principle** to account for quasi-specular reflection in a practical manner
- JCSDA has agreed to support (beginning Sep 2019) an upgrade to the CRTM IRSSE model as part of their 2019 Annual Operating Plan
 - The plan is to include surface temperature dependence along with some other misc upgrades
 - JPSS will support in-kind work until Sep 2019

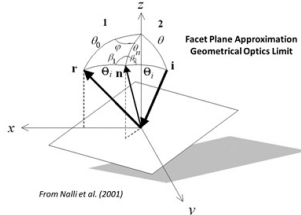
Radiative Transfer-Based Effective Emissivity (after Nalli et al. 2018a,b)

The **directional emissivity** of a terrestrial surface is defined as

$$\epsilon_{\nu}(\theta_0) \equiv \frac{I_{\nu s}(\theta_0)}{B_{\nu}(T_s)}$$

where the **surface-emitted radiance** (numerator) is separated from the **surface-leaving radiance** (as measured by a detector) by subtracting the **surface-reflected radiance**

$$I_{\nu s}(\theta_0) = R_{\nu s}(\theta_0, \varphi_0) - \frac{1}{\pi} \int_{\Omega} \int_{\theta} \int_{\varphi} \rho_{\nu}(\theta, \varphi; \theta_0, \varphi_0) I_{\nu}^{\downarrow}(\theta, \varphi) \cos(\theta) \sin(\theta) d\varphi d\theta$$



The **conical-directional reflectance** for non-isotropic incident radiation (Nicodemus et al. 1977) for the sea surface reflectance may be written as

$$\rho_{\nu}(\theta_0, \sigma_0^2) = \frac{\iint P(\theta_n, \varphi_n; \theta_0) P(\theta_n, \theta_0; \sigma_0^2) [B_{\nu}(T_s) - I_{\nu}^{\downarrow}(\theta)] d\varphi_n d\mu_n}{\iint P(\theta_n, \theta_0; \sigma_0^2) [B_{\nu}(T_s) - I_{\nu}^{\downarrow}(\theta)] d\varphi_n d\mu_n}$$

which, from the mean value theorem is equivalent to

$$\rho_{\nu}(\theta_0, \sigma_0^2) \equiv \rho_{\nu}(\bar{\theta}_n, \bar{\varphi}_n; \theta_0; \sigma_0^2) = \rho_{\nu}(\bar{\Theta}_e(\theta_0), \sigma_0^2)$$

The denominator simplifies as

$$\iint P(\theta_n, \theta_0; \sigma_0^2) [B_{\nu}(T_s) - I_{\nu}^{\downarrow}(\theta)] d\varphi_n d\mu_n = B_{\nu}(T_s) - I_{\nu}^{\downarrow}(\bar{\theta}_n)$$

where $\bar{\theta}_n \approx \theta_0$ is a diffusivity angle, thus allowing simplification of the surface-leaving radiance RTE as

$$R_{\nu s}(\theta_0) = B_{\nu}(T_s) - \rho_{\nu}(\bar{\Theta}_e, N_{\nu}) [B_{\nu}(T_s) - I_{\nu}^{\downarrow}(\bar{\theta}_n)] \approx B_{\nu}(T_s) - \rho_{\nu}(\Theta_e, N_{\nu}) [B_{\nu}(T_s) - I_{\nu}^{\downarrow}(\theta_0)]$$

Then, defining an **effective emissivity** as

$$\mathcal{E}_{\nu}(\theta_0) \equiv 1 - \rho_{\nu}[\Theta_e(\theta_0)],$$

where the mean incidence angle is the equivalent of an **effective emission angle**, $\bar{\Theta}_e \equiv \Theta_e$, one may arrive at a **simplified RTE for the SLR**

$$R_{\nu s}(\theta_0) = \mathcal{E}_{\nu}(\theta_0) B_{\nu}(T_s) + [1 - \mathcal{E}_{\nu}(\theta_0)] I_{\nu}^{\downarrow}(\theta_0)$$

The effective emissivity as defined is thus equivalent to

$$\mathcal{E}_{\nu}(\theta_0) = \frac{R_{\nu s}(\theta_0) - I_{\nu}^{\downarrow}(\theta_0)}{B_{\nu}(T_s) - I_{\nu}^{\downarrow}(\theta_0)}$$

The effective emission angle Θ_e is determined iteratively via **least-squares spectral variance minimization**

$$\sigma^2(\Delta\nu) = \frac{1}{n-1} \sum_{\nu} [T_{\nu s}(\Theta_e) - \bar{T}_{\nu s}(\Theta_e)]^2$$

where $T_{\nu s}(\Theta_e)$ is the skin temperature given by

$$T_{\nu s}(\Theta_e) = B_{\nu}^{-1} \left[\frac{R_{\nu s}(\theta_0) - \rho_{\nu}(\Theta_e, N_{\nu}) I_{\nu}^{\downarrow}(\theta_0)}{1 - \rho_{\nu}(\Theta_e, N_{\nu})} \right]$$

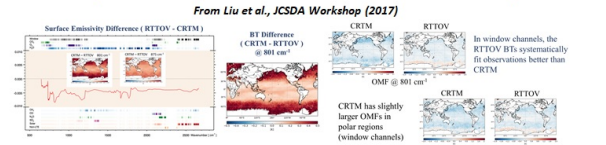
The retrieved Θ_e can then be used to derive the entire effective emissivity spectrum.

Need for Upgrades to the IRSSE Model

- CRTM IR effective-emissivity (IRSSE) model was derived based on high-accuracy surface-based FTS observations (MAERI), sound theoretical principles, and the need for practical implementation within the CRTM
 - Model was shown to have significantly better agreement with observations over the conventional models (e.g., Masuda 2006; Wu and Smith 1997)
 - However, small residual differences were still found to exist under certain conditions
 - For example, small residual positive biases were found in the SWIR windows, as well as LWIR under dry/cold atmospheric conditions
 - The author was aware of this at the time but ultimately favored simplicity (and was not supported for continued work)
- More importantly, however, was the realization of a **significant temperature dependence**
 - Nalli et al. (2008b): "In agreement with other recent work on the subject, we found a significant temperature dependence, which, if unaccounted for, can lead to spectral SLR errors of the same order of magnitude as those we have sought to correct. Therefore, additional work is desirable to derive an optimal seawater refractive index dataset..."
 - Unfortunately, this work was not supported in the time
 - However, the recent findings of (e.g., Liu et al. 2017) revealed significant **systematic bias** (as much as 1 K) on a **global scale**, thus bringing this issue back into focus for support

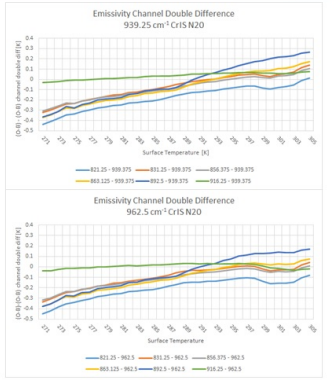
IRSSE Model Improvement

Temperature Dependence: Observed Global Scale Impact



Global OBS - CALC double-differences

- 2-weeks global NOAA-20 Cris data (OBS) versus CRTM model calculations (CALC)
- Shown are microwindow-channel double-differences of OBS - CALC in regions of varying surface temperature dependence observed in the IR spectrum (e.g., Pinkley et al. 1977; see the figures to the right)
- The double-differences serve to place control on the unknown atmospheric path uncertainties (e.g., model bias, cloud contamination, water vapor errors, etc.)
- Significant surface-temperature dependence is clearly visible on the order of 1 K
- This is of first order significance within the context of the total CRTM forward model uncertainty



Planned Work

- Temperature dependence
 - Refractive indices with temperature dependence are one possibility, however
 - Newman et al. (2005) is limited to a portion of the LWIR
 - Hale et al. (1972) is limited to 2 useful surface temperatures
 - Thus the tentative plan is to use data from Pinkley et al. (1977)
 - From the spectral variance minimization of $T_{\nu s}$ derive new LUT to include $T_{\nu s}$, i.e., $E(\nu, \theta_0, T_{\nu s})$
 - Validate using global data (e.g., Liu et al. 2017) as well as new MAERI campaigns-of-opportunity (e.g., Gero et al. 2016; Nalli et al. 2008b)
 - Collaboration with UW/CIMSS and UW/RSMMS
- Improvement to known residual biases in current model (viz., the SWIR window)
 - Perform separate spectral variance minimizations for the LWIR and SWIR bands (instead of the one-size-fits-all approach) and include $T_{\nu s}$ parameter for improved global fitting
 - Use latest version of FLRTM and use additional atmospheres in the training sample
- Doventail efforts with ongoing NUCAPS emissivity retrieval development and SARTA model upgrades



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